

DIELECTRICS -BOUNDARY VALUE PROBLEMS

# Flectricity and Magnetizm

Project PHYSNET Physics Bldg. Michigan State University East Lansing, MI

## DIELECTRICS - BOUNDARY VALUE PROBLEMS by R. D. Young

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#### Title: Dielectrics - Boundary Value Problems

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### Input Skills:

- 1. Vocabulary: dielectric, dielectric constant, permittivity, electric susceptibility, linear isotropic dielectric, polarization, electric displacement vector, external charge, polarization charge (MISN-0-506); Poisson's equation (MISN-0-505).
- 2. Express the solution to Laplace's equation in terms of zonal harmonics and cylindrical harmonics (MISN-0-505).
- 3. State Gauss's law for the electric displacement (MISN-0-506).

## Output Skills (Knowledge):

- K1. State the boundary conditions for the electric field and the displacement vector at an interface between two dielectric media.
- K2. State the forms of Poisson's and Laplace's equations for fields in the presence of dielectric material.

#### **Output Skills (Problem Solving):**

S1. Given a simple geometrical arrangement of two dielectric media or a dielectric medium in conjunction with conducting surfaces, use the boundary value conditions to determine the potential, electric field, displacement vector and charge densities in the media.

#### External Resources (Required):

 J. Reitz, F. Milford and R. Christy, Foundations of Electromagnetic Theory, 4th Edition, Addison-Wesley (1993).

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# by

## R. D. Young

## 1. Introduction

In general, two or more dielectric media are present in a given situation. In fact, the vacuum can always be considered as a dielectric with permittivity  $\epsilon_0$ . Or there may be at least one dielectric medium and one conducting medium in a problem. In these more complicated problems, the behavior of the field vectors  $\vec{E}$  and  $\vec{D}$  across the interface between two media is required for a solution. The boundary conditions across such an interface are derived in terms of the tangential component of the electric field  $\vec{E}$  and the normal component of the displacement  $\vec{D}$ . Poisson's equation in the presence of a dielectric medium is derived and specialized to Laplace's equation in the absence of any free charge.

Thus, an electrostatic problem involving linear, isotropic, and homogeneous dielectrics reduces to finding solutions of Laplace's equation in each medium, and joining the solutions in the various media by means of the boundary conditions alluded to above.

# 2. Procedures

- 1. Read Secs. 4-7 to 4-8 of the text.
- 2. Write down or mark in the text the condition on the tangential component of the electric field  $\vec{E}$  (eq. 4-42b) and the normal component of the displacement vector  $\vec{D}$  (eq. 4-41b) at an interface between two media. Also, write down a sentence or two explaining each of the equations (4-41b) and (4-42b). Be prepared to write down both the equations and the explanatory sentences on the Unit Test.
- 3. Write down or mark in the text Poisson's and Laplace's equations (4-48 and 4-49, respectively) in the presence of dielectric material.
- 4. Read Example 4-2, "Dielectric sphere in a uniform electric field," very carefully. This is a prototype example of solving Laplace's equation when a dielectric medium is present using boundary conditions on  $\vec{E}$  and  $\vec{D}$ .

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- 5. Read the *Supplementary Notes* for other examples of solving electrostatic problems involving linear, isotropic, and homogeneous dielectrics. The problems solved in the Notes are numbers 4-7, 4-16, and 4-17 of the text.
- 6. Solve the following problems:

4-8, 4-10, 4-15, 4-17

# 3. Supplementary Notes

1. Problem 4-7

Given:

Two dielectric media with dielectric constants  $K_1$  and  $K_2$ . Media separated by plane interface. No free charge on interface.

The angles  $\theta_1$  and  $\theta_2$  are the angles that the displacement vector makes with a normal to the interface in medium 1 and 2, respectively (see Fig. 1).

Find - Relationship between  $\theta_1$ ,  $\theta_2$  and  $K_1$ ,  $K_2$ .

Since there is no free charge,

$$D_{1n} = D_{2n} \tag{1}$$

$$E_{2t} = E_{1t} \tag{2}$$





where the subscripts  $\boldsymbol{n}$  and  $\boldsymbol{t}$  mean normal and tangential, respectively. Thus,

$$D_1 \cos \theta_1 = D_2 \cos \theta_2$$
$$E_1 \sin \theta_1 = E_2 \sin \theta_2$$
$$\frac{K_1}{K_2} = \frac{\tan \theta_1}{\tan \theta_2}$$



## 2. Problem 4-16

## Given:

Two parallel conducting plates - separated by d and maintained at potential difference  $\Delta V$ .

Slab (dielectric constant K) and uniform thickness d between plates. Slab does *not* completely fill volume between plates.

Find -

- a) Electric field in dielectric
- b) Electric field in vacuum
- c) Charge density  $\sigma$  on plate in contact with dielectric
- d) Charge density  $\sigma$  on plate in contact with vacuum
- e) Bound charge density  $\sigma_p$  on surface of dielectric

Solution -

In dielectric,

V = ax + b

In vacuum,

$$V = cx + \ell$$

But,

$$V(x = 0) = V_1$$
 and  $V(x = d) = V_2$ 

Then,

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$$V(x=0) = V_1 = b = \ell$$

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Also,

$$V(x = d) = V_2 = ad + V_1 = cd + V_1$$

Thus,

$$a = \frac{V_2 - V_1}{d} = \frac{\Delta V}{d}$$

. . .

Hence,

$$V(x) = \frac{\Delta V}{d}x + V_1$$

In vacuum and dielectric. Now,

$$E(x) = -\frac{\partial V}{\partial x} = -\frac{\Delta V}{d}$$

in vacuum and dielectric. Inside the conducting plates, D = E = 0. So, the boundary condition on the displacement vector (which has only a normal component) is

 $D = \sigma$ 

Thus,

$$\sigma = D = K\epsilon_0 E = -K\epsilon_0 \frac{\Delta V}{d}$$

on plate in contact with dielectric.

Thus,

$$\sigma = D = \epsilon_0 E = -\epsilon_0 \frac{\Delta V}{d}$$

on plate in contact with vacuum.

From  $\sigma_p = \vec{P} \cdot \hat{n}$ ,

$$\sigma_p = P = D - \epsilon_0 E = K \epsilon_0 E - \epsilon_0 E = (K - 1) \epsilon_0 E$$

So,

$$\sigma_p = -(K-1)\epsilon_0 \frac{\Delta V}{d}$$



Given:

Conducting sphere (radius R, free charge Q). Sphere floats halfsubmerged in a liquid dielectric of permittivity  $E_1$ . Region above sphere is gas of permittivity  $E_2$ .

Find - Electric field in dielectrics.

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Figure 3.

Solution -

Select a coordinate system with origin at the center of the sphere, the z-axis perpendicular to the surface of the liquid, and the x-y plane in the surface of the liquid. Then, the problem should be symmetric about the z-axis. Thus, the potential in medium 1 and 2 should be expanded in terms of zonal harmonics. So,

Medium 1

$$V_1 = A_1 + C_1 r^{-1} + A_2 r \cos \theta + C_2 r^{-2} \cos \theta + \frac{1}{2} A_3 r^2 (3 \cos^2 \theta - 1) + \frac{1}{2} C_3 r^{-3} (3 \cos^2 \theta - 1) + \dots$$

where

$$R < r < \infty, \frac{\pi}{2} \le \theta \le \pi$$

Medium 2

$$V_2 = A'_1 + C'_1 r^{-1} + A'_2 r \cos \theta + C'_2 r^{-2} \cos \theta + \frac{1}{2} A'_3 r^2 (3 \cos^2 \theta - 1) + \frac{1}{2} C'_3 r^{-3} (3 \cos^2 \theta - 1) + \dots$$

where

$$R < r < \infty, \frac{\pi}{2} \le \theta \le \pi$$

Let  $r \to \infty$ . Then,  $V_1$  and  $V_2$  must behave like a point charge potential. So,

$$A_2 = A_3 = \ldots = A'_2 = A'_3 = \ldots = 0$$

Likewise, as  $r \to R$ ,  $V_1$  and  $V_2$  must become constant independent of  $\theta$ . So,

$$C_2 = C_3 = \ldots = C'_2 = C'_3 = \ldots = 0$$

Hence,

$$V_1 = A_1 + C_1 r^{-1}$$
 and  $V_2 = A'_2 + C'_1 r^{-1}$ 

Thus,  $\vec{E}$  is a radial field. Since  $\vec{E}$  is independent of  $\theta$  and continuous at the interface,

$$\vec{E}_1 = \vec{E}_2 = E\frac{\vec{r}}{r}$$

The displacement vector D is parallel to E so D is radial. Apply Gauss' Law to an (imaginary) spherical surface of radius r, concentric with the conducting sphere. Then,

$$\oint \vec{D} \cdot \hat{n}dS = Q = D_1 \frac{A}{2} + D_2 \frac{A}{2}$$

where  $A = 4\pi r^2$ . So,

$$Q = \epsilon_1 E \cdot 2\pi r^2 + \epsilon_2 E \cdot 2\pi r^2$$
$$E = \frac{Q}{2\pi(\epsilon_1 + \epsilon_2)r^2}$$

and

$$\vec{E} = \frac{Q}{2\pi(\epsilon_1 + \epsilon_2)r^3}\vec{r}$$

You can calculate the charge densities.

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