

RESONANCES AND POLES; REAL AND IMAGINARY WORLDS
by
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## Input Skills:

1. Describe the time-average steady state power transferred into a damped driven oscillator from its driving force (MISN-0-31).
2. Plot pole trajectories of any given reciprocal of a quadratic function (MISN-0-59).

## Output Skills (Knowledge):

K1. Suppose there is a narrow resonance in a physical system and state what measurements you could make in order to determine the approximate locations of the nearby poles. State the conditions under which the approximate locations are accurate.

## Output Skills (Rule Application):

R1. Sketch complex-plane pole trajectories for given single functions of frequency.
R2. Given the observed width and position of a resonance, determine the approximate position of a nearby pole.

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## RESONANCES AND POLES; REAL AND IMAGINARY WORLDS

by
Peter Signell

## 1. Locating the Poles

It can be shown that the time-average steady-state power fed into a damped driven oscillator is: ${ }^{1}$

$$
\begin{equation*}
P_{\mathrm{ave}}(\omega)=\frac{F_{0}^{2} \omega_{0}^{2} \gamma / m}{\left(\omega_{0}^{2}-\omega^{2}\right)^{2}+4 \gamma^{2} \omega^{2}} \tag{1}
\end{equation*}
$$

where $\omega$ is the driving frequency (hence the frequency of oscillation of the oscillator), $F_{0}$ is the amplitude of the sinusoidal driving force, $\omega_{0}$ is the frequency of the oscillator when undamped and undriven, and $\gamma$ is the damping constant. If we consider the frequency of oscillation $\omega$ to be a complex variable, then the denominator can be factored (zeros found) by applying the quadratic root formula with $\omega^{2}$ as the variable. Then the roots of the denominator, $\omega_{p}$, are solutions to:

$$
\omega_{p}^{2}=-2 \gamma^{2}+\omega_{0}^{2} \pm 2 i \sqrt{\omega_{0}^{2}-\gamma^{2}}
$$

In turn, the square root of $\omega_{p}^{2}$ gives the actual roots:

$$
\begin{equation*}
\omega_{p}= \pm \sqrt{\omega_{0}^{2}-\gamma^{2}} \pm i \gamma \tag{2}
\end{equation*}
$$

$\triangleright$ Square this to prove that it is indeed the square root.
We can now write:

$$
\begin{equation*}
P_{\mathrm{ave}}(\omega)=\frac{F_{0}^{2} \omega^{2} \gamma / m}{\left(\omega-\omega_{1}\right)\left(\omega-\omega_{2}\right)\left(\omega-\omega_{3}\right)\left(\omega-\omega_{4}\right)} \tag{3}
\end{equation*}
$$

where the four $\omega$ 's are the four values one obtains with the four possible sign combinations in Eq. (2). At each of these roots of the denominator the value of $P$ becomes infinite so $P$ is said to have a simple pole there. The pole locations have an obvious symmetry (see Fig. 1). Note that the radius vector to any pole has the length:

$$
\sqrt{\left(\mathcal{R} e\left\{\omega_{p}\right\}\right)^{2}+\left(\mathcal{I} m\left\{\omega_{p}\right\}\right)^{2}}=\omega_{0}
$$

[^0]

Figure 1. Pole positions are shown by +'s.
which is independent of the amount of damping in the system. Thus as you increase the damping the poles trace out trajectories which are arcs of circles of constant radius $\omega_{0}$.
$\triangleright$ Plot these arcs on the graph above.
$\triangleright$ When you reach $\gamma=\omega_{0}$ the poles are all on the imaginary axis. Where do the trajectories go as you continue to increase $\gamma$ beyond $\omega_{0}$ ?
$\triangleright$ What happens as $\gamma \rightarrow \infty$ ?
$\triangleright$ How do these trajectories correlate with under damping, critical damping and overdamping? ${ }^{2}$ We now go back to the small-damping case, $\gamma \ll \omega_{0}$, where the poles in the first and fourth quadrants have the positions:

$$
\omega_{1,4}=\sqrt{\omega_{0}^{2}-\gamma^{2}} \pm \gamma \simeq \omega_{0} \pm i \gamma
$$

In this case the real parts of these pole positions are both very close to the resonant frequency $\omega_{0}$, as shown in Fig. 1.

$$
{ }^{2} \text { See "Damped Mechanical Oscillations" (MISN-0-29). }
$$



Figure 2. Definitions of resonance parameters.

## 2. Resonance Width

The width $\Gamma$ of the resonance at half-maximum can be deduced by writing:

$$
\begin{equation*}
P\left(\omega_{0} \pm \Gamma / 2\right)=\frac{1}{2} P\left(\omega_{0}\right) \tag{4}
\end{equation*}
$$

The resonance parameters are illustrated in Fig. 2.
For the case $\gamma \ll \omega_{0}$, the resonance will turn out to have a width $\Gamma$ which is very small compared to the resonant frequency $\omega_{0}$ so that the denominator of $P\left(\omega_{0}+\Gamma / 2\right)$ can be written:

$$
\left[\omega_{0}^{2}-\left(\omega_{0} \pm \Gamma / 2\right)^{2}\right]^{2}+4 \gamma^{2}\left(\omega_{0} \pm \Gamma / 2\right)^{2} \simeq \omega_{0}^{2}\left(\Gamma^{2}+4 \gamma^{2}\right)
$$

Putting this and $\omega^{2} \simeq \omega_{0}^{2}$ into (4) yields:

$$
\frac{F_{0}^{2} \omega_{0}^{2} \gamma / m}{\omega_{0}^{2}\left(\Gamma^{2}+4 \gamma^{2}\right)}=\frac{1}{2} \frac{F_{0}^{2} \omega_{0}^{2} \gamma / m}{4 \gamma^{2} \omega_{0}^{2}}
$$

for which the solution is $\gamma=\Gamma / 2$. Incidentally, this result confirms that for small damping, $\gamma \ll \omega_{0}$, we have a narrow width: $\Gamma \ll \omega_{0}$.

Then the imaginary parts of the pole positions for $\gamma \ll \omega_{0}$ are given by the half-width at half-maximum of the observed resonance. Thus as damping is made smaller ( $\gamma$ smaller), the poles approach the real axis from each side and the resonance gets narrower and higher.

The Appendix shows you the case: $\gamma=0.209 \omega_{0}$.
$\triangleright$ In the Appendix figures, color the $\operatorname{Im}\{\omega\}$ axis red, the vertical surface along the positive $\mathcal{R} e\{\omega\}$ axis blue.

## Acknowledgments

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Figure 3. Definitions of the three axes for Figures 4 and 5.

## A. Pictures of Poles and Resonances



Figure 4. $P_{\text {ave }}(\omega)$. The surface has been removed for $\mathcal{I} m\{\omega\}>0$.

## PROBLEM SUPPLEMENT

These problems also occur on this module's Model Exam.

1. Locate the poles of: $F(\omega)=\frac{a^{2}}{\left(\omega-\omega_{0}\right)^{2}+a^{2}}$
2. Sketch the pole trajectories resulting from variations of the parameter $a$ in problem (1).
3. For the above case, determine the relations between the pole positions and the resonance width and position.
4. If there is a narrow resonance in a physical system, state what measurements you could make in order to determine the approximate locations of the nearby poles. State the conditions under which the approximate locations are accurate.

## Brief Answers:

1. Solve the denominator for $\omega$.
2. Your plot should show the two poles fleeing the real axis in opposite directions along a single straight line as $a$ is increased.
3. You should find, good for all values of $a: \Gamma=2 \operatorname{I} m\left\{\omega_{p}\right\}$, and $\omega_{\text {res }}=$ $\mathcal{R} e\left\{\omega_{p}\right\}$.
4. See this module's text, and think about it.

## MODEL EXAM

These problems also occur in this module's Problem Supplement.

1. Locate the poles of: $F(\omega)=\frac{a^{2}}{\left(\omega-\omega_{0}\right)^{2}+a^{2}}$
2. Sketch the pole trajectories resulting from variations of the parameter $a$ in problem (1).
3. For the above case, determine the relations between the pole positions and the resonance width and position.
4. If there is a narrow resonance in a physical system, state what measurements you could make in order to determine the approximate locations of the nearby poles. State the conditions under which the approximate locations are accurate.

[^0]:    ${ }^{1}$ See "Damped Driven Oscillations; Mechanical Resonances" (MISN-O-31).

