

COMPUTER ALGORITHM FOR THE DAMPED DRIVEN OSCILLATOR


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> Peter Signell

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## Title: Computer Algorithm for the Damped Driven Oscillator

Author: Peter Signell, Dept. of Physics, Mich. State Univ
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## Input Skills:

1. Expand a given function about a given point using Taylor's series (MISN-0-4).

## Output Skills (Knowledge):

K1. Given the force acting on a damped driven oscillator along with the oscillator's position and velocity at a specified time, derive a Numerov type algorithm for the approximate numerical calculation of the oscillator's position at all past and future times.

## Post-Options:

1. "Response of a Damped Driven Oscillator" (MISN-0-30).
2. "Damped Driven Oscillations; Mechanical Resonances" (MISN-031).
3. "Laplace Transform Solution for the Damped Driven Oscillator" (MISN-0-47).

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Our publications are designed: (i) to be updated quickly in response to field tests and new scientific developments; (ii) to be used in both classroom and professional settings; (iii) to show the prerequisite dependencies existing among the various chunks of physics knowledge and skill, as a guide both to mental organization and to use of the materials; and (iv) to be adapted quickly to specific user needs ranging from single-skill instruction to complete custom textbooks.

New authors, reviewers and field testers are welcome.

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| Andrew Schnepp | Webmaster |
| :--- | :--- |
| Eugene Kales | Graphics |
| Peter Signell | Project Director |

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$$

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by<br>\section*{Peter Signell}

## 1. Derivation

1a. Equation for a Damped Driven Oscillator. The equation to be solved is that of the damped harmonically-driven oscillator ${ }^{1}$ acted on by the force:

$$
F(t)=-k x(t)-\lambda v(t)+F_{0} \cos (\omega t)
$$

Since ${ }^{2}$

$$
F(t)=m a(t)=m \frac{d x(t)}{d t} \equiv m x^{\prime \prime}(t)
$$

and

$$
v(t)=\frac{d x(t)}{d t} \equiv x^{\prime}(t)
$$

our equation to be solved is:

$$
\begin{equation*}
m x^{\prime \prime}(t)+k x(t)+\lambda x^{\prime}(t)-f(t)=0 \tag{1}
\end{equation*}
$$

where

$$
f(t) \equiv F_{0} \cos (\omega t)
$$

1b. Approximate Derivatives by Finite Differences. We now make two power series expansions ${ }^{3}$ about time t:

$$
\begin{aligned}
& x(t+\Delta)=x(t)+\Delta x^{\prime}(t)+\left(\Delta^{2} / 2\right) x^{\prime \prime}(t)+\left(\Delta^{3} / 6\right) x^{\prime \prime \prime}(t)+\ldots \\
& x(t-\Delta)=x(t)-\Delta x^{\prime}(t)+\left(\Delta^{2} / 2\right) x^{\prime \prime}(t)-\left(\Delta^{3} / 6\right) x^{\prime \prime \prime}(t)+\ldots
\end{aligned}
$$

We will choose $\Delta$ sufficiently small so that we can disregard all terms after $\Delta^{2}$ without incurring much error. Then add and subtract the above equations to obtain:

$$
x(t+\Delta)+x(t-\Delta) \simeq 2 x(t)+\Delta^{2} x^{\prime \prime}(t)
$$

[^0]

Figure 1. The labeling system for time net-points.

$$
x(t+\Delta)-x(t-\Delta) \simeq 2 \Delta x^{\prime}(t)
$$

or

$$
\begin{gathered}
x^{\prime \prime}(t) \simeq \frac{1}{\Delta^{2}}[x(t+\Delta)-2 x(t)+x(t-\Delta)] \\
x^{\prime}(t) \simeq \frac{1}{2 \Delta}[x(t+\Delta)-x(t-\Delta)]
\end{gathered}
$$

which are often quoted in calculus courses. ${ }^{4}$
1c. Net-Point Times. We need a more succinct labeling system at this point or we won't be able to see the forest for the trees. We define discrete "net-point" times as: $t_{n} \equiv n \Delta$, where $\Delta$ is some small time interval (see Fig. 1). We can then rewrite the derivatives at time $t_{n}$ as:

$$
\begin{gather*}
x_{n}^{\prime \prime} \simeq \frac{1}{\Delta^{2}}\left(x_{n+1}-2 x_{n}+x_{n-1}\right)  \tag{2}\\
x_{n}^{\prime} \simeq \frac{1}{2 \Delta}\left(x_{n+1}-x_{n-1}\right) \tag{3}
\end{gather*}
$$

Then at time $t_{n}$ our basic Eq. (1) becomes:

$$
m x_{n}^{\prime \prime}+k x_{n}+\lambda x_{n}^{\prime}-f_{n}=0
$$

Substituting for $x_{n}^{\prime \prime}$ and $x_{n}^{\prime}$ :

$$
\frac{m}{\Delta^{2}}\left(x_{n+1}-2 x_{n}+x_{n-1}\right)+k x_{n}+\frac{\lambda}{2 \Delta}\left(x_{n+1}-x_{n-1}\right)-f_{n}=0
$$

[^1]Collecting coefficients of the $x$ values, we get:

$$
\begin{equation*}
x_{n+1}=A x_{n}+B x_{n-1}+\Delta^{2} C_{n} \tag{4}
\end{equation*}
$$

where:

$$
\begin{aligned}
A & =\left(2-\Delta^{2} \frac{k}{m}\right) / D \\
B & =-\left(1-\Delta \frac{\lambda}{2 m}\right) / D \\
C_{n} & =\left(\frac{f_{n}}{m}\right) / D \\
D & =1+\Delta \frac{\lambda}{2 m}
\end{aligned}
$$

## 2. Application

2a. Iterating the Recurrence Relation. Equation (4) is called a three point recurrence relation because it relates $x$ at three successive time net-points. If you know $x_{0}$ and $x_{1}$, you can use this relation to calculate $x_{2}$. Then knowing $x_{1}$ and $x_{2}$, you can calculate $x_{3}$, etc. You can keep up this recurrence procedure until you have reached the time for which you wish to know the solution. For example, this might be a time sufficiently large so that transient effects have died away and only the steady state solution remains.


Figure 2. The Calculated Curve. Draw a smooth curve through the plotted points and you have $x(t)$.

2b. Starting the Iteration. With a three-point recurrence relation one can calculate values at successive time net-points, but one needs two adjacent values in order to get started. In the example cited in the preceding paragraph, one needed $x_{0}$ and $x_{1}$ in order to start the procedure. Those two initial values generally come from a specification of a "complete set of initial conditions." Our Eq. (1), being a second order equation for $x$, always requires that one specify two independent conditions on $x$.

A common case is where both the position and velocity are known at some instant of time. Here we will call that time $t=0$; hence we know $x_{0}$ and $x_{1}$ by solving Eq. (3) at $n=0$ for $x_{-1}$ (Note: $v_{0}=x_{0}^{\prime}$ ), and using that to eliminate $x_{-1}$ from Eq. (4) for $n=0$. We thus obtain $x_{1}$ in terms of $x_{0}$ and $v_{0}$ (usually given quantities):

$$
x_{1}=\left(1-\Delta^{2} \frac{k}{2 m}\right) x_{0}+\left(\Delta-\Delta^{2} \frac{\lambda}{2 m}\right) v_{0}+\Delta^{2} \frac{f_{0}}{2 m}
$$

Now the algorithm is complete. Put in $x_{0}$ and $v_{0}$ to get $x_{0}$ and $x_{1}, x_{0}$ and $x_{1}$ to get $x_{2}, x_{1}$ and $x_{2}$ to get $x_{3}$, etc.
$\triangleright$ How would you obtain position values for times earlier than the one for which $x$ and $v$ are initially known? Answer: To see the answer, replace each of the following letters by its successor in the alphabet: mdfzshud cdksz.

## Acknowledgments

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## MODEL EXAM

1. The force acting on a damped driven oscillator is $F=-k x-\lambda v+$ $F_{0} \cos (\omega t)$; the oscillator's initial position and speed are $x_{0}$ and $v_{0}$. Derive without notes, a complete Numerov-type algorithm for the (approximate) calculation of the oscillator's position at all times. NOTE: The algorithm is to include the method of use of any three point recurrence relation such as:

$$
x_{n+1}=A x_{n}+B x_{n-1}+\Delta^{2} C_{n}
$$

where:

$$
\begin{aligned}
A & =\left(2-\Delta^{2} \frac{k}{m}\right) / D \\
B & =-\left(1-\Delta \frac{\lambda}{2 m}\right) / D \\
C_{n} & =\left(\frac{f_{n}}{m}\right) / D \\
D & =1+\Delta \frac{\lambda}{2 m}
\end{aligned}
$$

and boundary conditions such as:

$$
x_{1}=\left(1-\Delta^{2} \frac{k}{2 m}\right) x_{0}+\left(\Delta-\Delta^{2} \frac{\lambda}{2 m}\right) v_{0}+\Delta^{2} \frac{f_{0}}{2 m}
$$

## Brief Answers:

1. See this module's text.

[^0]:    ${ }^{1}$ See "Damped Driven Oscillations; Mechanical Resonances" (MISN-0-31).
    ${ }^{2}$ See "Particle Dynamics - The Laws of Motion" (MISN-0-14),
    ${ }^{3}$ See "Taylor's Series for the Expansion of a Function About a Point" (MISN-0-4).

[^1]:    ${ }^{4}$ The second equation above is used in "Response of a Damped Driven Oscillator" (MISN-0-30) where a microcomputer is utilized in finding solutions.

